Optically-active Vibrations and Elastic Constants of Calcite and Aragonite¹⁾

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Treatments of the normal coordinate and the elastic constant were undertaken for two calcium carbonate crystals. By the use of the least-squares technique, we obtained sets of the force constants which reproduce satisfactorily both the observed vibrational frequencies and the macroscopic elastic constants. The intermolecular interactions, the elastic strains of the internal coordinates, and the nature of the splitting of the normal frequencies in these crystals were briefly discussed.

The vibrational spectra of molecular crystals and ionic crystals have been studied by many investigators. Some of them have performed normal coordinate analyses, and the force fields of many crystals have been determined. Recently more rigorous treatments of the lattice dynamics, such as those based on the rigid-ion model²) or on the shell model,³) have been made for relatively simple crystal systems mainly by physicists.⁴) Presumably the impracticability of these treatments arises in part from the difficulties in the reasonable estimation of the potential parameters. The results of these treatments are quite suggestive for the analysis of the crystal vibration in a complex crystal system, though the application of the above models to actual complex crystals is difficult at the present time.

Generally, in contrast to the treatment of isolated free molecules, many restrictions which are characteristic of the solid state make it difficult to determine the force fields from only the vibrational spectra. These restrictions are caused by the experimental limitations, the complicated crystal structures, the strong intra- or inter-molecular interactions, and so on. The short-range interaction, intramolecular or intermolecular, is considered to dominate the crystal vibration and the elastic deformation, at least in molecular crystals or in predominantly covalent crystals. Therefore, it is necessary, for the determination of the force fields in the successional models, to ascertain the fundamental short-range force fields.

The present investigations of the short-range force fields of two calcium carbonates, and of sodium nitrate and sodium nitrite in the low-temperature ferroelectric phase¹⁾ were undertaken not only in order to check the reasonability of the vibrational force fields, but also to determine those force fields which interpret both the observed vibrational frequencies and the elastic constants.

Methods of Calculation

Crystal Structure. According to the X-ray diffraction analyses, 5,6) the Bravais unit cells of calcite and aragonite belong to the rhombohedral space group $D_{3d}^6-R\bar{3}c$, with Z=2, and the orthorhombic $D_{2h}^{16}-Pmcn$, with Z=4 respectively. The cell constants are $a_{\rm rh}=6.3753$ Å and $\alpha=46^{\circ}4.6'$ in calcite and a=4.9614 Å, b=7.9671 Å, and c=5.7404 Å in aragonite. Figures 1 and 2 show the structures of calcite and aragonite respectively. The structure of calcite may be regarded as a sodium chloride arrangement

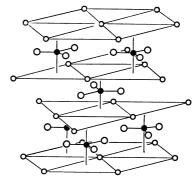


Fig. 1. Crystal structure of calcite.

•: carbon, \bigcirc : calcium, \bigcirc : oxygen.

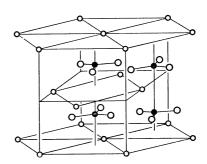


Fig. 2. Crystal structure of aragonite.

•: carbon, \bigcirc : calcium, \bigcirc : oxygen.

distorted by the introduction of disk-shaped carbonate groups. Each calcium atom is coordinated by six oxygen atoms, each belonging to a different carbonate group, and such oxygen atom is bound to two calcium atoms. On the other hand, the orthorhombic structure of aragonite is related to that of nickel arsenide. Each calcium atom is now surrounded by nine neighbors, each of which is, in its turn, bound to three calcium atoms.

Normal Coordinate Treatment. The application of the GF-matrix method⁷⁾ to the analysis of optically-active crystal vibrations was reported previously by Shimanouchi, Tsuboi, and Miyazawa.⁸⁾ The modified Urey-Bradley force field⁹⁾ was applied to the intramolecular vibrations in carbonate groups, the correction terms of which were the bond interaction term, p, for the adjacent C–O bond pairs and the angle interaction term, n, for the adjacent O–C–O angle pairs in a single carbonate group. On the other hand, the central force field, which is based on the two-body interaction, was used for the intermolecular vibrations.

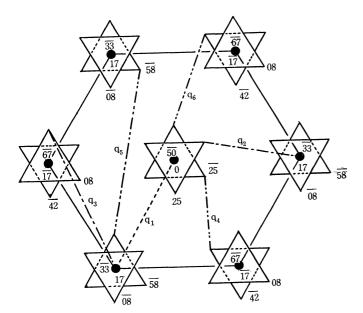


Fig. 3. Short-range interactions in calcite viewed along the z axis of the hexagonal cell. The heights of atoms or carbonate groups are expressed in units of $100/c_0$.

----, calcium···carbon; —----, calcium···oxygen; —----, oxygen···oxygen.

In this study, only those pairs that lie within 3.5 Å were taken into consideration. As is illustrated in Figs. 3 and 4, there are six different kinds of these pairs in calcite and fifteen in aragonite respectively. The corresponding atomic separations are also listed in Table 3.

In many cases, group theoretical considerations, such as the site-group analysis¹⁰ or the factor-group analysis,¹¹ have been found to be useful for the analysis

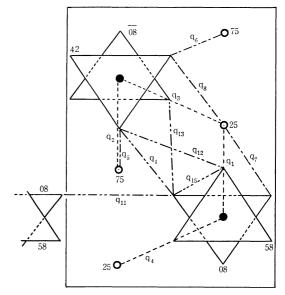


Fig. 4. Short-range interactions in aragonite viewed along the z axis. Other descriptions are the same to Fig. 3.

of crystal vibrations. The factor-group analysis of these crystals was once performed by Bhagavantam and Venkatarayudu;¹¹⁾ their results are summarized in Tables 1 and 2. It is obvious from these tables that, of the twenty-seven degrees of freedom for genuine vibrations eight modes are infrared-active, and five, Raman-active, in calcite, and that, of the fifty-seven degrees, twenty-seven are infrared-active, and thirty, Raman-active, in aragonite. The selection rules for the optically-active vibrations are followed under the mutual exclusion rule, since each crystal has a center of symmetry.

Table 1. Factor-group analysis of calcite

| D_{3d}^6 | Е | $2S_6$ | $2S_6^2$ | $S_6^{\ 3}$ | $3C_2$ | 3 σ _v | N_i | T_i | T_i' | R_{i}' | n_i | IR | Raman |
|---------------------------|---|--------|------------|-------------|--------|------------------|-------|-------|--------|----------|-------|----------------|---|
| A _{1g} | 1 | 1 | 1 | 1 | 1 | 1 | 1 | 0 | 0 | 0 | 1 | f | $\alpha_{xx} = \alpha_{yy}, \alpha_{zz}$ |
| A_{1u} | 1 | -1 | 1 | -1 | 1 | -1 | 2 | 0 | 1 | 0 | 1 | f | f |
| ${ m A_{2g}}$ | 1 | 1 | 1 | 1 | -1 | — 1 | 3 | 0 | 1 | 1 | 1 | \mathbf{f} | f |
| A_{2u} | 1 | -1 | 1 | -1 | -1 | 1 | 4 | 1 | 1 | 1 | 1 | $\mathbf{M_z}$ | \mathbf{f} |
| $\mathbf{E}_{\mathbf{g}}$ | 2 | -1 | -1 | 2 | 0 | 0 | 4 | 0 | 1 | 1 | 2 | \mathbf{f} | $\alpha_{xx} = -\alpha_{yy}, \alpha_{xy}, \alpha_{yz}, \alpha_{zx}$ |
| $\mathbf{E_u}$ | 2 | 1 | — 1 | -2 | 0 | 0 | 6 | 1 | 2 | 1 | 2 | M_x , M_y | f |

f, forbidden; N_i , number of degrees of freedom belonging to the *i*-th species; T_i , number of over-all translations of Bravais unit cell; T_i' , number of translational lattice modes; R_i' , number of rotational lattice modes; n_i , number of intramolecular vibrational modes.

TABLE 2. FACTOR-GROUP ANALYSIS OF ARAGONITE

| $\mathrm{D_{2h}^{16}}$ | E | $\mathbf{C_2^c}$ | C_2^b | C_2^a | i | $\sigma_{ m c}$ | $\sigma_{ m b}$ | $\sigma_{\rm a}$ | N_i | T_i | T_{i}' | R_i' | n_i | IR | Raman |
|----------------------------|---|------------------|------------|---------|------------|-----------------|-----------------|------------------|-------|-------|----------|--------|-------|---------------------------|---|
| A_{1g} | 1 | 1 | 1 | 1 | 1 | 1 | 1 | 1 | 9 | 0 | 4 | 1 | 4 | f | $\alpha_{xx}, \alpha_{yy}, \alpha_{zz}$ |
| A_{1u} | 1 | 1 | 1 | 1 | -1 | -1 | -1 | -1 | 6 | 0 | 2 | 2 | 2 | f | \mathbf{f} |
| $ m B_{1g}$ | 1 | 1 | 1 | 1 | 1 | 1 | -1 | — 1 | 6 | 0 | 2 | 2 | 2 | f | α_{xy} |
| $\mathrm{B_{1u}}$ | 1 | 1 | — 1 | -1 | -1 | -1 | 1 | 1 | 9 | 1 | 3 | 1 | 4 | \mathbf{M}_{z} | f |
| $\mathrm{B_{2g}}$ | 1 | -1 | 1 | -1 | 1 | -1 | 1 | 1 | 6 | 0 | 2 | 2 | 2 | f | α_{zx} |
| ${f B_{2u}}$ | 1 | -1 | 1 | 1 | -1 | 1 | -1 | 1 | 9 | 1 | 3 | 1 | 4 | $\mathbf{M}_{\mathbf{y}}$ | \mathbf{f} |
| $\mathrm{B}_{3\mathbf{g}}$ | 1 | 1 | — 1 | 1 | 1 | -1 | -1 | 1 | 9 | 0 | 4 | 1 | 4 | \mathbf{f} | $lpha_{	ext{yz}}$ |
| $\mathrm{B_{3u}}$ | 1 | -1 | -1 | 1 | — 1 | 1 | 1 | -1 | 6 | 1 | 1 | 2 | 2 | $\mathbf{M}_{\mathbf{x}}$ | f |

Notations are the same to Table 1.

Table 3. Force constants (mdyn/Å) of calcite and aragonite

| Ca | lcite | | A | ragonite | : |
|--|-------|------------------------|--|----------|------------------------|
| Force cons | stant | Distance ^{a)} | Force con | stant | Distance ^{a)} |
| $\overline{K(\text{C-O})}$ | 5.520 | 1.290 | K(C-O) | 5.669 | 1.282 |
| $H(OCO)^{b)}$ | 0.340 | | $H(OCO)^{b)}$ | 0.340 | |
| F(O O)b) | 1.720 | | H(O O)b) | 1.720 | |
| $\alpha(\mathrm{CO_3})^{\mathrm{c}_{\!1}}$ | 0.797 | | $\alpha(\mathrm{CO_3})^{\mathrm{c}_3}$ | 0.738 | |
| þ | 0.052 | | þ | -0.125 | |
| n^{c} | 0.021 | | n^{c} | 0.016 | |
| $f_1(Ca \ C)$ | 0.103 | 3.185 | $f_1(Ca \ C)$ | 0.288 | 2.904 |
| | | | $f_2(\text{Ca C})$ | 0.074 | 2.937 |
| | | | $f_3(\text{Ca C})$ | 0.030 | 3.249 |
| | | | $f_4(\text{Ca C})$ | 0.020 | 3.415 |
| $f_2(\text{Ca O})$ | 0.425 | 2.333 | $f_5(\text{Ca O})$ | 0.251 | 2.419 |
| | | | $f_6(\text{Ca O})$ | 0.096 | 2.445 |
| | | | $f_7(\text{Ca O})$ | 0.040 | 2.550 |
| | | | $f_8(\text{Ca O})$ | 0.060 | 2.520 |
| $f_3(\mathbf{C} \mathbf{O})^{\mathrm{b}}$ | 0.000 | 3.392 | $f_9(\text{Ca O})$ | 0.040 | 2.653 |
| | | | $f_{10}(C C)$ | 0.123 | 2.877 |
| | | | $f_{11}(O \ O)$ | 0.316 | 2.744 |
| | | | $f_{12}(O \ O)$ | 0.224 | 3.065 |
| $f_4(O O)$ | 0.080 | 3.185 | $f_{13}(O O)$ | 0.132 | 2.983 |
| $f_5(O O)$ | 0.040 | 3.203 | $f_{14}(O \ O)$ | 0.086 | 3.075 |
| $f_6(O O)$ | 0.028 | 3.392 | $f_{15}(O \ O)$ | 0.069 | |

- a) Unit of distances is in Å. b) Fixed force constants.
- c) Units of n and α are in mdyn·Å.

Treatment of the Elastic Constant. When a crystal is subjected to homogeneous deformation, the elastic constant of the crystal can be calculated in accordance with the method established by Shiro. The course of the calculation is outlined below. In the first place, by the use of the resulting matrices of the preceding section, i.e., F_0^R and B_ρ , and of the external strain matrix, D, the (6×6) elastic constant matrix is derived as:

$$\boldsymbol{C} = 1/v \tilde{\boldsymbol{H}} \boldsymbol{F_0}^{\mathrm{R}} \boldsymbol{H}, \tag{1}$$

where a tilde denotes a transposed matrix, where v is the unit cell volume, and where

$$\boldsymbol{H} = (\boldsymbol{E} - \boldsymbol{B}_{\rho} (\tilde{\boldsymbol{B}}_{\rho} \boldsymbol{F}_{0}^{R} \boldsymbol{B}_{\rho})^{-1} \tilde{\boldsymbol{B}}_{\rho} \boldsymbol{F}_{0}^{R}) \boldsymbol{D}. \tag{2}$$

The Jacobian matrix with respect to the elastic constant, $(\partial C_i/\partial K_h)$, and the contribution of the *h*-thp otential term to the *i*-th elastic constant, $(PED)_{ih}$, are also given:

$$J_{ih}^{E} = (\partial C_{i}/\partial K_{h}) = 1/v(\tilde{\boldsymbol{H}}\boldsymbol{A}_{h}\boldsymbol{H})_{ii}, \tag{3}$$

$$(PED)_{ih} = J_{ih}^{E} K_h / C_i, \tag{4}$$

respectively, where A_h is the coefficient matrix of the k-th term in the potential energy matrix. Furthermore, the change in the internal coordinates due to the elastic stress, ΔS is calculated as:

$$\Delta \mathbf{R}_0 = \mathbf{H} \mathbf{C}^{-1} \Delta \mathbf{S}. \tag{5}$$

The change, $\Delta \mathbf{R}_0$, is found to be useful for the analysis of the elastic strains.

Refinement of the Force Constant. The Jacobians with respect to the optically-active vibrations, J^{λ} , and to the elastic constants, J^{E} , are derived independently

from the treatments of the preceding sections. The force constants of the force fields were refined by the least-squares method:

$$\mathbf{\Delta K} = (\tilde{\mathbf{J}}^{\lambda} \mathbf{W}_{\lambda} \mathbf{J}^{\lambda} + \tilde{\mathbf{J}}^{E} \mathbf{W}_{E} \mathbf{J}^{E})^{-1}
\times (\tilde{\mathbf{J}}^{\lambda} \mathbf{W}_{\lambda} \Delta \lambda + \tilde{\mathbf{J}}^{E} \mathbf{W}_{E} \Delta \mathbf{C}),$$
(6)

where W^{λ} and W^{E} are diagonal weight matrices for the observed vibrational frequencies and elastic constants respectively, where $\Delta \lambda$ is a deviation column matrix of the calculated frequencies from the corresponding observed ones, and where ΔC is that of the elastic constants.

Results and Discussion

Optically-active Vibrations. The vibrational spectra of calcite and aragonite have been reported by several authors. 13-19) However, it seems likely that the observed resonance frequencies will deviate slightly from the corresponding eigenfrequencies²⁰⁾ when one come to think of the ionic properties of the two crystals. In our normal coordinate treatments, we postulated an equilibrium crystal field, i.e., any dynamical effects are frozen out, since only the short-range interactions were taken into consideration. Consequently, the eigenfrequencies in the idealized equilibrium field should generally deviate from those in the actual dynamical lattices. We have, however, regarded the observed resonance frequencies as the eigenfrequencies without any correction, because the appropriate estimation of the shifts is impracticable, especially for a complex crystal system. Provided that the shifts are brought about through interactions among the vibrating dipoles, the observed frequencies of the g-symmetric vibrations are found to be approximately equal to those in the equilibrium crystal field.²¹⁾ Therefore, we have given some priority to the g-symmetric vibrations over the u-symmetric ones in the course of the refinement of the force constants in calcite.

Ramdas'¹³) and Tsuboi's data¹⁴) were adopted as the observed frequencies for the intramolecular vibrations of the carbonate groups and for the (intermolecular) lattice vibrations respectively. In the case of aragonite, though few data are available, Couture's data,¹³) which were obtained from the depolarization study of Raman scattering, are noteworthy. Hence, only Couture's data were taken as the observed data, the remaining data^{16–19}) other than Couture's amounted to no more than supplimentary references.

Intramolecular Vibration of the Carbonate Groups: Carbonate groups in calcite occupy a site, the symmetry of which is $\mathrm{D_3}$, while it is $\mathrm{C_s}$ in aragonite. Because of the lowering of the site symmetry, more numbers of the optically-active frequencies are to be observed in aragonite than in calcite. The intramolecular vibrations in these crystals may, in principle, be explained in terms of the force field of an isolated free carbonate group, since the observed intramolecular frequencies of these crystals show much similarity. The results on force constants are shown in Table 3. The observed and calculated frequencies in calcite and aragonite are also compared in Tables 4 and 5 respectively. It

Table 4. Observed and calculated frequencies (cm⁻¹) of calcite

| Sym. | | Obsd.a) | Calcd. | Assignment | Sym. | | Obsd.a) | Calcd. | Assignment |
|---------------------------|------------|---------|--------|-----------------------------------|---------------------------|------------|---------|--------|-----------------------------------|
| A _{1g} | ν_{11} | 1086 | 1086 | CO str. | A _{1u} | v_{41} | | 1089 | CO str. |
| _ | | | | | | v_{42} | | 211 | Ca trans(z). |
| A_{2g} | v_{21} | | 890 | CO_3 out-of-plane def. | | | | | |
| | v_{22} | | 287 | CO_3 trans(z). | A_{2u} | v_{51} | 885 | 884 | CO3 out-of-plane def |
| | $ u_{23} $ | | 197 | $CO_3 \operatorname{rot}(z)$. | | v_{52} | 290 | 334 | Ca, CO_3 trans(z). |
| | | | | | | v_{53} | 110 | 120 | $CO_3 \operatorname{rot}(z)$. |
| $\mathbf{E}_{\mathbf{g}}$ | v_{31} | 1434 | 1428 | CO str. | | | | | |
| - | $ u_{32} $ | 712 | 710 | OCO def. | $\mathbf{E}_{\mathbf{u}}$ | v_{61} | 1484 | 1428 | CO str. |
| | ν_{33} | 284 | 266 | $CO_3 \operatorname{rot}(x, y)$. | | v_{62} | 706 | 711 | OCO def. |
| | v_{34} | 156 | 169 | CO_3 trans(x, y). | | ν_{63} | 320 | 312 | $CO_3 \operatorname{rot}(x, y)$. |
| | | | | | | v_{64} | 230 | 224 | Ca $trans(x, y)$. |
| | | | | | | v_{65} | 110 | 113 | Ca, CO_3 trans(x, y). |

a) See Refs. 13 and 14.

Table 5. Observed and calculated frequencies (cm^{-1}) of aragonite

| Sym. | | Obsd.a) | Calcd. | Assignment | Sym. | | Obsd.b | Calcd. | Assignment |
|----------------------------|------------|---------|--------|-----------------------------------|-------------------|------------|--------|--------|----------------------------------|
| A _{1g} | ν_{11} | 1464 | 1468 | CO str. | A_{1u} | v_{51} | | 1456 | CO str. |
| -0 | v_{12} | 1086 | 1090 | CO str. | | v_{52} | | 685 | OCO def. |
| | ν_{13} | 854 | 853 | CO_3 out-of-plane def. | | v_{53} | | 222 | CO_3 rot(y). |
| | v_{14} | 707 | 717 | OCO def. | | v_{54} | | 140 | Ca, CO_3 trans(x). |
| | v_{15} | 287 | 275 | $CO_3 \operatorname{rot}(x)$. | | v_{55} | | 105 | $CO_3 \operatorname{rot}(z)$. |
| | v_{16} | 191 | 208 | CO_3 trans(y). | | v_{56} | | 99 | Ca trans(x). |
| | ν_{17} | 143 | 146 | Ca, CO_3 trans(z). | | | | | |
| | v_{18} | | 121 | Ca trans(y). | $\mathbf{B_{1u}}$ | ν_{61} | | 1468 | CO str. |
| | ν_{19} | 105 | 99 | Ca $trans(z)$. | | v_{62} | 1084 | 1079 | CO str. |
| | | | | | | v_{63} | 857 | 851 | CO ₃ out-of-plane def |
| B_{1g} | ν_{21} | 1415 | 1457 | CO str. | | v_{64} | 712 | 708 | OCO def. |
| *8 | v_{22} | 716 | 717 | OCO def. | | ν_{65} | 294 | 254 | $CO_3 \operatorname{rot}(x)$. |
| | v_{23} | | 248 | CO_3 rot(y). | | v_{66} | 200 | 185 | Ca, CO_3 trans(y). |
| | v_{24} | | 197 | CO_3 rot(z). | | v_{67} | | 152 | Ca trans (y) , (z) . |
| | v_{25} | 154 | 138 | Ca, CO ₃ trans(x). | | ν_{68} | 114 | 131 | Ca trans(y). |
| | v_{26} | 114 | 97 | Ca trans(y). | | | | | |
| | | | | | $\mathbf{B_{2u}}$ | v_{71} | 1480 | 1467 | CO str. |
| $\mathrm{B}_{2\mathbf{g}}$ | v_{31} | 1415 | 1457 | CO str. | | v_{72} | 1082 | 1084 | CO str. |
| -0 | v_{32} | 702 | 711 | OCO def. | | ν_{73} | 877 | 890 | CO ₃ out-of-plane def |
| | v_{33} | 263 | 247 | CO_3 rot(z). | | v_{74} | 706 | 711 | OCO def. |
| | v_{34} | 250 | 162 | CO_3 rot(y). | | v_{75} | 274 | 272 | Ca, CO_3 trans(z). |
| | v_{35} | 208 | 147 | Ca, CO_3 trans(x). | | v_{76} | 200 | 208 | $CO_3 \operatorname{rot}(x)$. |
| | v_{36} | 183 | 73 | Ca, CO_3 trans(x). | | v_{77} | | 169 | Ca, CO_3 trans(y). |
| | | | | | | v_{78} | 100 | 87 | Ca, CO_3 trans(z). |
| $\mathbf{B_{3g}}$ | v_{41} | 1466 | 1470 | CO str. | | | | | |
| | v_{42} | 1086 | 1084 | CO str. | $\mathbf{B_{3u}}$ | ν_{81} | 1480 | 1456 | CO str. |
| | v_{43} | 854 | 890 | CO ₃ out-of-plane def. | | v_{82} | 712 | 691 | OCO def. |
| | v_{44} | 703 | 718 | OCO def. | | v_{83} | 274 | 163 | CO_3 rot(y). |
| | v_{45} | 276 | 273 | CO_3 trans(z). | | ν_{84} | 200 | 137 | Ca, CO_3 trans(x). |
| | v_{46} | 222 | 235 | CO_3 trans(y). | | v_{85} | 118 | 123 | $CO_3 \operatorname{rot}(z)$. |
| | v_{47} | 193 | 201 | CO_3 rot(x). | | | | | |
| | v_{48} | 113 | 128 | Ca trans(y). | | | | | |
| | v_{49} | 105 | 117 | Ca trans(z). | | | | | |

a) See Ref. 15. b) See Ref. 11.

is obvious from these tables that the large splittings of the frequencies between g-symmetric and u-symmetric C-O stretching vibrations cannot be explained on the basis of the short-range interaction model. Not only the coupling of the motion of the carbonate groups, but also the interaction between the vibrating dipoles would cause these large splittings. These consideration have been made by Haas and Ketelaar, ²¹⁾ Decius, ²²⁾ and Ishigame et al. ²³⁾ Their results are quite suggestive, but they give only qualitative explanations for complicated crystal systems. These models will actually be applied in our subsequent papers.

For the observed Raman frequency shifts at 1415 cm⁻¹ and around 1460 cm⁻¹ in aragonite, there are two ways of interpretation. Krishnan¹⁶) regarded only the shift at 1463 cm⁻¹ as the fundamental band and assigned the shift at 1415 cm⁻¹ to the overtone band of the fundamental band at 707 cm⁻¹ on the basis of his second-order Raman scattering study. On the other hand, Krishnamurti¹⁷) considered all these shifts to be fundamental bands. In this study, we have regarded the shift at 1415 cm⁻¹ as the fundamental band, with reference to the spectra of carbonate or nitrate crystals.

Lattice Vibration: Normal coordinate treatments of calcite have been performed by Venkateswarlu and Ramaswamy,²⁴⁾ Tsuboi,¹⁴⁾ Shiro,²⁵⁾ and Plihal and Schaack;²⁶⁾ however, no treatment of aragonite has been reported. Though only the short-range interactions were taken into consideration in both Tsuboi's¹⁴⁾ and Shiro's²⁵⁾ treatments, their results are satisfactory for at least the factor-group vibrations. Assuming an idealized simpler unit cell which contains only four ions, Plihal and Schaack²⁶⁾ took the long-range Coulomb interaction of vibrating dipoles and the influence of the electronic polarizability of the component ions as well as

the short-range interaction; thus they explained fairly well the observed TO-LO splittings of the optically-active lattice vibrations in the calcite-form crystals. A detailed analysis based on the rigid ion model is now in progress; the results will be published in time.

Therefore, we have determined the short-range fields, which reproduce both the observed vibrational frequencies and the elastic constants, and have also improved the previous treatment.²⁵⁾ Tables 4 and 5 show the observed and calculated frequencies and the tentative assignments of the observed fundamental bands. The contributions of all the potential terms to the calculated vibrational frequencies are also given in Tables 6 and 7. The results of the calculation imply that certain lattice modes in the same symmetry species are fairly well coupled with each other, but that the couplings between an intramolecular mode and lattice modes are negligibly smal, except for the out-of-plane angle deformation modes in carbonate groups. As is shown in Tables 4 and 5, the calculated frequencies of aragonite agree less well with the observed ones than that of calcite, particularly in the $B_{2\mathbf{g}}$ and $B_{\epsilon\mathbf{u}}$ symmetric vibrations. Our resulting force field for calcite is substantially that obtained by Tsuboi. 14) The interatomic force constants will be discussed briefly in a later section.

Elastic Constants. The elastic constants of these crystals have previously been determined ^{27,28)} and tabulated in the "Landolt-Börnstein Phys. Chem. Tables." However, only a few attempts at theoretical calculation have yet been reported. For example, calculations of the two particular constants were performed by Oke^{30,31)} on the basis of the electrostatic theory of crystal lattices. We have ourselves calculated the elastic constants by the method described above; the calculated elastic constants are compared

Table 6. Contribution (%) of potential terms to calculated frequencies for calcite

| Mode | K | Н | \overline{F} | α | þ | n | f_1 | f_2 | f_3 | f_4 | f_5 | f_6 |
|------------|----|----|----------------|----|------------|-----------|-------|------------|-------|-------|-------|-------|
| ν_{11} | 50 | 0 | 46 | 0 | 1 | 0 | 0 | 2 | 0 | 0 | 1 | 0 |
| v_{41} | 50 | 0 | 46 | 0 | 1 | 0 | 0 | 2 | 0 | 0 | 1 | 0 |
| v_{42} | 1 | 0 | 1 | 0 | 0 | 0 | 12 | 86 | 0 | 0 | 0 | 0 |
| v_{21} | 0 | 0 | 0 | 96 | 0 | 0 | 2 | 1 | 0 | 1 | 0 | 0 |
| v_{22} | 0 | 0 | 0 | 1 | 0 | 0 | 3 | 70 | 0 | 4 | 8 | 14 |
| $ u_{23} $ | 0 | 0 | 0 | 0 | 0 | 0 | 4 | 4 | 0 | 63 | 23 | 6 |
| $ u_{51} $ | 0 | 0 | 0 | 97 | 0 | 0 | 2 | 1 | 0 | 1 | 0 | 0 |
| $ u_{52} $ | 0 | 0 | 0 | 0 | 0 | 0 | 5 | 88 | 0 | 0 | 5 | 1 |
| v_{53} | 0 | 0 | 0 | 0 | 0 | 0 | 22 | 1 | 0 | 0 | 67 | 9 |
| ν_{31} | 85 | 9 | 6 | 0 | 1 | 0 | 1 | 0 | 0 | 0 | 0 | 0 |
| v_{32} | 3 | 30 | 59 | 0 | 0 | -1 | 1 | 6 | 0 | 1 | 2 | 0 |
| v_{33} | 0 | 0 | 1 | 0 | 0 | 0 | 2 | 64 | 0 | 25 | 0 | 8 |
| v_{34} | 0 | 0 | 0 | 0 | 0 | 0 | 21 | 32 | 0 | 44 | 0 | 4 |
| v_{61} | 85 | 9 | 6 | 0 | — 1 | 0 | 1 | 0 | 0 | 0 | 0 | 0 |
| v_{62} | 3 | 29 | 58 | 0 | 0 | -1 | 1 | 7 | 0 | 0 | 2 | 1 |
| v_{63} | 0 | 1 | 1 | 0 | 0 | 0 | 13 | 81 | 0 | 0 | 0 | 4 |
| v_{64} | 0 | 1 | 3 | 0 | 0 | 0 | 21 | 7 3 | 0 | 0 | 0 | 2 |
| $ u_{64} $ | 0 | 0 | 0 | 0 | 0 | 0 | 43 | 1 | 0 | 0 | 0 | 55 |

Notations are the same to Tables 3 and 4.

Table 7. Contribution (%) of potential terms to calculated frequencies for aragonite

| Mode | K | H | F | α | þ | n | f_1 | f_2 | f_3 | f_4 | f_5 | f_{6} | f_7 | f_8 | f_9 | f_{10} | f_{11} | f_{12} | f_{13} | f_{14} | f_{15} |
|--|---------------|---------|---------|-----------|----------------|-----------|----------|---------------|----------|----------|----------|-----------|---------|--------|---------------|----------|---------------|----------|----------|----------------|----------|
| v_{11} | 83 | 8 | 5 | 0 | 2 | 0 | 1 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 |
| ν_{12} | 51 | 0 | 46 | 0 | -2 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 3 | 1 | 0 | 1 | 0 |
| ν_{13} | 0 | 0 | 0 | 98 | 0 | 0 | 1 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 |
| v_{14} | 3 | 29 | 58 | 0 | 0 | -1 | 1 | 0 | 0 | 0 | 0 | 0 | 0 | 1 | 0 | 0 | 6 | 2 | 0 | 0 | 0 |
| v_{15} | 0 | 0 | 0 | 0 | 0 | 0 | 4 | 0 | 1 | 0 | 30 | 5 | 0 | 2 | 0 | 0 | 0 | 13 | 11 | 2 | 32 |
| $ u_{16} $ | 0 | 0 | 0 | 0 | 0 | 0 | 27 20 | 5 0 | 0 0 | 0 3 | 6 2 | $0 \\ 24$ | 7 6 | 7 0 | 1 1 | 0 | 0 | 24 13 | 0 5 | $\frac{3}{24}$ | 19 1 |
| $\begin{matrix}v_{17}\\v_{18}\end{matrix}$ | 0 | 0 | 0 | 0 | 0 | 0 | 20 | 0 | 11 | 3 | 32 | 4 | 0 | 3 | 2 | 0 | 0 | 0 | 0 | 2 | 41 |
| v_{19} | 0 | 0 | 0 | 0 | 0 | 0 | 8 | 7 | 1 | 1 | 3 | 19 | 3 | 10 | 3 | 0 | 0 | 28 | 8 | 2 | 7 |
| | 84 | | | | | | | | | | | | | | | | | | | 0 | 0 |
| v_{21} | 3 | 8 30 | 5 58 | 0 | 2 | $0 \\ -1$ | 0 | 1 0 | 0 | 0 | 0 0 | 0 | 0 | 0 | 0 1 | 0 0 | 0 | $0 \\ 4$ | 0 3 | 2 | 0 |
| $egin{array}{c} {v_{22}} \\ {v_{23}} \end{array}$ | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 9 | 0 | 4 | 2 | 0 | 0 | 28 | 14 | 4 | 39 |
| v_{24} | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 3 | 0 | 0 | 0 | 3 | 3 | 0 | 9 | 0 | 0 | 10 | 34 | 2 | 34 |
| v_{25} | 0 | 0 | 1 | 0 | 0 | 0 | 0 | 29 | 0 | 0 | 0 | 1 | 15 | 23 | 7 | 0 | 0 | 9 | 1 | 0 | 14 |
| v_{26} | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 8 | 0 | 15 | 0 | 35 | 0 | 0 | 11 | 0 | 0 | 7 | 9 | 1 | 14 |
| v_{31} | 84 | 8 | 5 | 0 | 2 | 0 | 0 | 1 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 |
| v_{32} | 3 | 30 | 59 | 0 | 0 | -1 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 1 | 0 | 0 | 4 | 3 | 0 | 0 |
| $ u_{33}^{52} $ | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 2 | 0 | 0 | 0 | 2 | 2 | 3 | 8 | 0 | 0 | 31 | 34 | 8 | 10 |
| v_{34} | 0 | 0 | 1 | 0 | 0 | 0 | 0 | 2 | 0 | 0 | 0 | 12 | 2 | 6 | 1 | 0 | 0 | 9 | 3 | 22 | 41 |
| $ u_{35} $ | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 26 | 0 | 5 | 0 | 26 | 2 | 18 | 8 | 0 | 0 | 5 | 1 | 0 | 9 |
| v_{36} | 0 | 1 | 1 | 0 | 0 | 0 | 0 | 5 | 0 | 1 | 0 | 0 | 38 | 0 | 12 | 0 | 0 | 2 | 25 | 6 | 9 |
| ν_{41} | 83 | 8 | 5 | 0 | 2 | 0 | 1 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 |
| v_{42} | 51 | 0 | 47 | 0 | -2 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 3 | 1 | 0 | 0 | 0 |
| v_{43} | $0 \\ 3$ | 0 29 | 0 58 | 90 0 | 0 | $0 \\ -1$ | 0 1 | 0 | $0 \\ 0$ | 0 0 | $0 \\ 0$ | 0 | 0 0 | 0 1 | $0 \\ 0$ | 7 0 | 0 6 | $0 \\ 2$ | 0 0 | 0 0 | 1 1 |
| $egin{array}{c} v_{44} \ v_{45} \end{array}$ | 0 | 0 | 0 | 0 | o | 0 | 4 | 1 | 1 | 1 | 22 | 4 | 1 | 0 | 0 | 15 | 0 | 0 | 6 | 0 | 44 |
| v_{46} | 0 | 0 | 0 | 0 | 0 | 0 | 12 | 1 | 1 | 0 | 12 | 4 | 2 | 3 | 0 | 2 | 0 | 42 | 7 | 2 | 11 |
| ν_{47} | 0 | 0 | 0 | 0 | 0 | 0 | 13 | 0 | 1 | 0 | 10 | 2 | 5 | 0 | 1 | 1 | 0 | 2 | 2 | 38 | 24 |
| v_{48} | 0 | 0 | 0 | 0 | 0 | 0 | 32 | 12 | 0 | 0 | 0 | 1 | 7 | 24 | 0 | 2 | 0 | 5 | 1 | 8 | 6 |
| v_{49} | 0 | | 0 | 0 | 0 | 0 | 0 | 1 | 4 | 4 | 24 | 36 | 0 | 0 | 5 | 3 | 0 | 1 | 6 | 0 | 14 |
| v_{51} | 84 3 | 8 32 | 5 63 | 0 | 2 0 | $0 \\ -1$ | 0 0 | 1 0 | 0 | 0 | 0 | 0 0 | 0 | 0 | 0 1 | 0 0 | 0 0 | 0 0 | $0 \\ 0$ | 0 0 | 0 0 |
| $rac{{v_5}_2}{{v_5}_3}$ | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 13 | 2 | 2 | 0 | 0 | 0 | 0 | 0 | 9 | 73 |
| v_{54} | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 27 | 0 | 0 | 0 | 3 | 0 | 1 | 31 | 0 | 0 | 0 | 0 | 21 | 16 |
| v_{55} | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 12 | 0 | 0 | 0 | 1 | 13 | 51 | 2 | 0 | 0 | 0 | 0 | 17 | 2 |
| v_{56} | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 5 | 0 | 14 | 0 | 30 | 16 | 0 | 9 | 0 | 0 | 0 | 0 | 8 | 18 |
| $ u_{61} $ | 83 | 8 | 5 | 0 | 2 | 0 | 1 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 |
| v_{62} | 52 0 | 0 | 47 0 | 0 · 98 | $-2 \\ 0$ | 0 | 0 1 | 0 0 | 0 | 0 | 0 0 | 0 0 | 0 0 | 0 | 0 | $0 \\ 0$ | $\frac{3}{0}$ | 0 0 | $0 \\ 0$ | 0 0 | 0 0 |
| $v_{63} \\ v_{64}$ | 3 | 30 | 59 | 90 | | -1 | 1 | 0 | 0 | 0 | 0 | 0 | 0 | 1 | 0 | 0 | 5 | 0 | 0 | 0 | 0 |
| v_{65} | 0 | 0 | 0 | ő | 0 | Ô | ō | o | o | o | 31 | o | o | 0 | 2 | Ö | 0 | 0 | 0 | 15 | 51 |
| v_{66} | 0 | 0 | 0 | 0 | 0 | 0 | 42 | 0 | 2 | 0 | 18 | 2 | 11 | 0 | 1 | 0 | 0 | 0 | 0 | 16 | 7 |
| ν ₆₇ | 0 | 0 | 0 | 0 | 0 | 0 | 13 | 1 | 4 | 5 | 18 | 42 | 1 | 1 | 1 | 1 | 0 | 0 | 0 | 0 | 14 |
| ν_{68} | 0 | 0 | 0 | 0 | 0 | 0 | 11 | 16 | 3 | 0 | 4 | 3 | 3 | 34 | 0 | 0 | 0 | 0 | 0 | 17 | 8 |
| v_{71} | 83 51 | 8 0 | 5 46 | 0 0 - | $\frac{2}{-2}$ | 0 | 1 0 | 0 0 | 0 0 | $0 \\ 0$ | 0 0 | 0 0 | 0 0 | 0 0 | 0 0 | 0 | $0 \\ 3$ | 0 0 | $0 \\ 0$ | 0 1 | 0 0 |
| $egin{array}{c} u_{72} \\ u_{73} onumber \end{array}$ | 0 | 0 | 0 | 90 | - 2 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 7 | 0 | 0 | 0 | 0 | 1 |
| v_{74} | 3 | 30 | 59 | 0 | | 1 | 1 | ő | Ö | 0 | o 0 | 0 | ő | 1 | ő | ó | 5 | 0 | 0 | 0 | 1 |
| v_{75} | 0 | 0 | 0 | 0 | 0 | 0 | 3 | 0 | 1 | 0 | 19 | 1 | 1 | 0 | 1 | 16 | 0 | 0 | 0 | 10 | 48 |
| v_{76} | 0 | 0 | 0 | 0 | 0 | 0 | 11 | 2 | 1 | 2 | 33 | 18 | 1 | 7 | 1 | 3 | 0 | 0 | 0 | 0 | 20 |
| v_{77} | $\frac{1}{0}$ | 0 0 | 0 | 0 | 0 | 0 | 43 0 | $\frac{6}{3}$ | 1 7 | 0 6 | 9 4 | 4 16 | 13 0 | 9 6 | $\frac{2}{0}$ | 0 0 | $0 \\ 0$ | 0 0 | 0 0 | 1 3 | 12 53 |
| $ u_{78} $ | | | | | | | | | | | | | | | | | | | | | |
| v_{81} | 84 3 | 8 32 | 5 62 | $0 \\ 0$ | 2 | 0 -1 | 0 0 | $\frac{1}{0}$ | $0 \\ 0$ | $0 \\ 0$ | $0 \\ 0$ | $0 \\ 0$ | 0 0 | 0 0 | 0 1 | 0 0 | 0 | 0 0 | $0 \\ 0$ | $0 \\ 2$ | 0 1 |
| $egin{array}{c} { u_{82}} \\ { u_{83}} \end{array}$ | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 12 | 0 | 2 | 0 | 38 | 6 | 0 | 9 | 0 | 0 | 0 | 0 | 7 | 25 |
| v_{84} | ő | 0 | 0 | 0 | 0 | 0 | 0 | 17 | o | 3 | Ö | 0 | 5 | 34 | 3 | 0 | 0 | 0 | 0 | 7 | 30 |
| v_{85} | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 6 | 0 | 1 | 0 | 0 | 11 | 4 | 31 | 0 | 0 | 0 | 0 | 1 | 46 |

Notations are the same to Tables 3 and 5.

Table 8. Observed and calculated elastic constants (1010 dyn/cm2) of calcite and aragonite

| T) 6 | | Calcite | | D 16 | Aragonite | | | |
|-----------------------|---------|---------|---|------------------------|-----------|--------|-------------------|--|
| $\mathrm{D_{3d}}^{6}$ | Obsd.a) | Calcd. | Sym. | $\mathrm{D_{2h}^{16}}$ | Obsd.b) | Calcd. | Sym. | |
| C_{11} | 144.0 | 118.7 | A _{1g} , E _g | C_{11} | 160.0 | 165.2 | A_{1g} | |
| C_{21} | 53.9 | 49.9 | A_{1g}, E_{g} | C_{21} | 37.0 | 39.1 | A_{1g} | |
| C_{31} | 51.1 | 52.8 | A_{1g} | C_{22} | 89.2 | 73.4 | A_{1g} | |
| C_{33} | 84.0 | 75.5 | A_{1g} | C_{31} | 1.7 | 19.8 | A_{1g} | |
| C_{41} | -20.5 | -20.9 | $\mathbf{E}_{\mathbf{g}}$ | C_{32} | 15.4 | 24.6 | A_{1g} | |
| C_{44} | 33.5 | 33.0 | $\mathbf{E}_{\mathbf{g}}^{^{\mathrm{c}}}$ | C_{33} | 84.2 | 89.2 | A_{1g} | |
| | | | | C_{44} | 41.2 | 23.2 | $\mathrm{B_{3g}}$ | |
| | | | | C_{55} | 25.6 | 17.3 | $\mathrm{B_{2g}}$ | |
| | | | | C_{66} | 42.7 | 24.1 | $\mathbf{B_{1g}}$ | |

a) See Ref. 27. b) See Ref. 28.

Table 9. Contribution (%) of potential terms to calculated elastic constants

| | | | | 4 | Aragoni | te | | | | | | | Calci | te | | |
|-----------|----|----|----|----|---------|----|----|----|----|--------------------|----|----|------------|----|----|----|
| | 11 | 21 | 22 | 31 | 32 | 33 | 44 | 55 | 66 | | 11 | 21 | 31 | 33 | 41 | 44 |
| K | 2 | 3 | 1 | 1 | 1 | 0 | 0 | 0 | 0 | K | 1 | 2 | 1 | 0 | 1 | 0 |
| H | 1 | -1 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | H | 0 | -1 | 0 | 0 | 2 | 1 |
| ${\it F}$ | 4 | 1 | 1 | 0 | 0 | 0 | 0 | 0 | 0 | F | 2 | 0 | 1 | 0 | 4 | 1 |
| α | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | α | 0 | 0 | 0 | 0 | 0 | 0 |
| þ | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | þ | 0 | 0 | 0 | 0 | 0 | 0 |
| n | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | n | 0 | 0 | 0 | 0 | 0 | 0 |
| f_1 | 1 | 5 | 4 | 2 | 2 | 0 | 7 | 0 | 0 | $f_{\mathfrak{t}}$ | 21 | 15 | 16 | 5 | 38 | 26 |
| f_2 | 5 | 13 | 4 | 12 | 6 | 1 | 4 | 7 | 0 | | | | | | | |
| f_3 | 0 | 2 | 2 | 1 | 3 | 0 | 5 | 0 | 0 | | | | | | | |
| f_4 | 1 | 3 | 1 | 6 | 3 | 1 | 3 | 9 | 2 | | | | | | | |
| f_5 | 0 | 1 | 1 | 5 | 8 | 4 | 4 | 0 | 0 | f_2 | 43 | 75 | 69 | 40 | 43 | 43 |
| f_{6} | 2 | 11 | 7 | 27 | 25 | 8 | 5 | 25 | 4 | | | | | | | |
| f_7 | 0 | 2 | 1 | 1 | 1 | 0 | 1 | 1 | 10 | | | | | | | |
| f_8 | 1 | 8 | 13 | 5 | 12 | 1 | 10 | 2 | 1 | | | | | | | |
| f_{9} | 5 | -3 | 0 | 6 | 1 | 0 | 0 | 6 | 2 | f_3 | 0 | 0 | 0 | 0 | 0 | 0 |
| f_{10} | 0 | 0 | 0 | 0 | -5 | 20 | 0 | 0 | 0 | | | | | | | |
| f_{11} | 69 | -2 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | | | | | | | |
| f_{12} | 0 | 4 | 8 | -1 | -3 | 0 | 10 | 0 | 16 | | | | | | | |
| f_{13} | 5 | 29 | 20 | 2 | 2 | 0 | 0 | 7 | 17 | f_4 | 2 | -5 | 5 | 33 | 20 | 23 |
| f_{14} | 2 | 23 | 28 | 28 | 52 | 9 | 8 | 13 | 37 | f_5 | 26 | 18 | — 1 | 0 | -2 | 0 |
| f_{15} | 0 | 0 | 7 | 4 | -6 | 56 | 42 | 30 | 5 | f_{6} | 4 | -4 | 9 | 20 | -5 | 6 |

Notations are the same to Tables 3 and 8.

with the observed ones in Table 8. It is obvious from these results that a good agreement was obtained for calcite, while only a rough fit was obtained for aragonite. Particularly, the agreement of the elastic constants related to the shear strains, i.e., C_{44} , C_{55} , and C_{66} , is relatively poor in aragonite because of the lack of effective oblique interaction terms compared with those in calcite. The contributions of all the potential terms to the calculated elastic constants are also listed in Table 9. This table shows that the elastic constants have a strong dependence on the following potential terms; f_1 , f_2 , f_4 , and f_5 in calcite, and f_2 , f_6 , f_8 , f_{13} , f_{14} , and f_{15} in aragonite. It is likely that these contributions are partly dependent on the spatial orientation of the internal coordinates. These contributions are considered to be independent of those to the vibrational frequencies, as is shown in Tables 6 and 7; therefore

the determination of the interatomic force constants becomes more meaningful.

Interatomic Force Constants. The resulting interatomic force constants are plotted as functions of the interatomic distance in Fig. 5. As this figure illustrates, three steep curves can be drawn through the points of the same pairs, and the distances of their burst upbreaks approximately correspond to the sums of the van der Waals radii or of the Pauling's ionic radii³²⁾ of the constituent atoms, as shown in Fig. 5: $R_{\text{Ca}} + R_{\text{O}} = 2.39 \,\text{Å}, R_{\text{Ca}} + R_{\text{C}} = 2.69 \,\text{Å}, \text{ and } 2R_{\text{O}} = 2.80 \,\text{Å}.$ Since these interatomic interaction force constants should be derived from the appropriate interatomic potentials, we began to estimate them with simple Lennard-Jones 6-9 type potentials³³⁾ determined for some purely ionic crystals and noble gas crystals; however, we thus obtained unreasonable results except

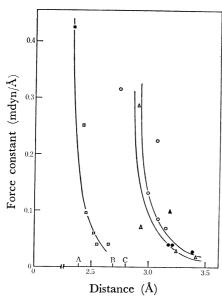


Fig. 5. Interatomic force constants. Closed and open figures indicate those for calcite and aragonite respectively.

 \triangle , calcium···carbon; \square , calcium···oxygen; \bigcirc , oxygen ···oxygen; $R_{\rm C}$, van der Waals radius of carbon; $R_{\rm O}$, van der Waals radius of oxygen; $R_{\rm Ca}$, Pauling's ionic radius of calcium.

A: $R_{Ca}+R_{O}$, B: $R_{Ca}+R_{C}$, C: R_{O}

for a few cases. We inevitably chose several trial sets of them which reproduced the observed values as well as possible, and then refined by the least-squares method.¹²⁾ As is shown in Fig. 5, some plots were found to deviate from the corresponding curves. This must mainly be because of the uncertainty of the experimental frequencies and/or elastic constants, and the incompleteness of the force fields. The former is considered to have an appreciable influence on the force constants in aragonite, while both the observed frequencies and elastic constants were found to be fairly reliable in calcite. The results listed in Tables 6 and 9 suggest that the force constant, f_1 , of calcite contributes particularly to certain lattice modes and to all the elastic constants except C_{33} . If this force constant is reduced to the value on the curve, less agreement will be obtained. Therefore, the large deviation or the closed triangle must arise mainly from the incompleteness of the force field. The incompleteness of the force field can be improved by the rigid ion treatment or by the polarizable ion treatment; these treatments are now in progress. As to the open circles in Fig. 5, f_{11} was found to contribute in particular to C_{11} (see Table 9), and f_{12} , to almost all the lattice modes in aragonite. Therefore, the incompleteness must exert an appreciable influence on the force constants, above all, the force constants which are closely related to the lattice vibrations in aragonite.

The force constants obtained would be useful as the initial values of the force constants in the normal coordinate treatments of similar crystals, though it is meaningless to discuss certain fundamental properties of the solid state on the basis of only these data.

Analysis of Elastic Strain. Table 10 shows the

Table 10. Change $(\mathring{A}^3/\text{mdyn})$ of internal coordinates due to elastic stress in calcite

| Coord.b) | S_{xx} | S_{yy} | S_{zz} | $S_{ m yz}$ |
|--|----------|----------|----------|-------------|
| CO str. | 1 | 1 | 0 | 1 |
| OCO def.a) | -3 | 3 | 0 | -2 |
| CO ₃ out-of-plane ^{a)} | 0 | 0 | 0 | 0 |
| q_1 | 4 | 42 | -1 | -70 |
| q_2 | 10 | 16 | 19 | 58 |
| q_3 | 27 | -65 | 121 | -156 |
| q_4 | -18 | -46 | 144 | -109 |
| q_5 | 142 | -26 | -72 | -110 |
| q_6 | -36 | -74 | 121 | 64 |

a) Unit of angle coordinates is in Rad \cdot Å ²/mdyn. b) Notations are the same to Fig. 3.

Table 11. Change (ų/mdyn) of internal coordinates due to elastic stress in aragonite

| Coord.b) | S_{xx} | S_{yy} | S_{zz} | S_{yz} | S_{zx} | S_{xy} |
|---------------------------------------|----------|-------------------|----------|----------------------------|----------------------------|----------------------------|
| | A_{1g} | $\mathbf{A_{1g}}$ | A_{1g} | $\mathrm{B}_{3\mathbf{g}}$ | $\mathrm{B}_{2\mathbf{g}}$ | $\mathrm{B}_{1\mathrm{g}}$ |
| CO str. | 2 | 0 | 0 | — 1 | 0 | -1 |
| | 0 | 1 | 0 | 0 | 0 | 0 |
| OCO def.a) | -4 | 2 | 0 | -1 | -1 | -9 |
| | 5 | -3 | 0 | 1 | 0 | 0 |
| CO_3 out-of- plane ^{a)} | 1 | 1 | 0 | 0 | 0 | 0 |
| q_1 | 4 | 15 | 1 | -27 | 0 | 0 |
| q_2 | 28 | 21 | 12 | -41 | 78 | 5 |
| q_3 | -2 | 42 | 7 | 73 | 0 | 0 |
| q_4 | 22 | 17 | 32 | 63 | -164 | -55 |
| q_5 | 2 | 7 | 19 | 23 | 0 | 0 |
| q_6 | 12 | 27 | 44 | 41 | -124 | -35 |
| q_7 | 9 | 16 | 6 | -26 | -40 | 88 |
| q_8 | -4 | 70 | 9 | -73 | 41 | -20 |
| q_9 | 51 | -37 | 19 | 2 | 95 | 41 |
| q_{10} | 3 | 18 | 70 | -1 | 0 | 0 |
| q_{11} | 63 | -28 | 4 | 0 | 0 | 0 |
| q_{12} | -3 | 31 | -10 | -37 | 4 | 47 |
| q_{13} | 12 | 54 | -8 | -7 | 57 | 62 |
| q_{14} | 1 | 77 | 40 | -54 | -95 | 114 |
| q_{15} | 10 | -55 | 91 | 52 | 104 | 23 |
| | 11 | -32 | 93 | -54 | -122 | 38 |
| | -6 | 19 | 89 | -116 | 5 | -14 |
| | | | | | | |

- a) Unit of angle coordinates is in Rad·Å²/mdyn.
- b) Notations are the same to Fig. 4.

change in the internal coordinates due to unit elastic stress in calcite. We can see from this table that the q_5 coordinate is sensitive to the $S_{\rm xx}$ applied stress, the q_3 , q_5 , and q_6 coordinates, to the $S_{\rm zz}$ stress, and the q_3 , q_4 , and q_5 coordinates, to the $S_{\rm yz}$ stress. The internal coordinates of calcite, therefore, can be changed by the applied stress along the direction z. For aragonite, Table 11 suggests that the q_4 , q_6 , q_{14} , and q_{15} coordinates a strong dependence on the appropriate stresses. As well as in calcite, it seems that calcium atoms in aragonite would be easy to displace in the directions to the adjacent carbonate groups, and that the carbonate would be susceptible to certain deformations along the

z axis. From these considerations, some information will be gained about the phase transitions between polymorphic crystals, though we cannot reach any definite conclusions on the basis of these data only.

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